

# TIME-DEPENDENT MODEL CALCULATIONS FOR $H_2^+$ IN A STRONG ULTRA-SHORT LASER PULSE

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## OVERVIEW:

- I. **Time Propagation** of the Schrödinger Equation
- II. **Visualization**
- III. **Extracting Physics**
- IV. **Conclusions and Outlook**

## Model for $\text{H}_2^+$ in a Strong Laser Field

- The basic idea was presented by Kulander *et al.*, Phys. Rev. A 53 (1996) 2562.
- Many other groups have worked on this and similar problems.
- A common simplification of the full problem is the **reduction to one dimension with two degrees of freedom**, one for the **nuclear separation (R)** and one for the **electronic motion along the internuclear (z) axis**, with the electric field also along this axis.
- **Goals of the present project :**
  - Apply our method, developed for e-H collisional excitation and ionization, to an explicitly time-dependent problem
  - Check a different time-propagation scheme (**leap-frog**)
  - **Visualization of the results**
  - **Extract physics by projection techniques**
  - Try to **avoid masking potentials** that reduce the norm of the wavefunction

- Going **beyond the Born-Oppenheimer approximation**, we solve the time-dependent Schrödinger Equation

$$i\frac{\partial\Psi(R, z, t)}{\partial t} = \left[ -\frac{1}{2\mu}\frac{\partial^2}{\partial R^2} - \frac{1}{2}\frac{\partial^2}{\partial z^2} + \frac{1}{\sqrt{R^2 + q_n}} - \frac{1}{\sqrt{(z - R/2)^2 + q_e}} - \frac{1}{\sqrt{(z + R/2)^2 + q_e}} + z f(t) \mathcal{E}_0 \sin(\omega t) \right] \Psi(R, z, t)$$

by time-propagating the initial state

$$\Psi(R, z, t = 0) = F_v^{1\sigma_g}(R) G_R^{1\sigma_g}(z)$$

- Here  $G_R^{1\sigma_g}(z)$  is the electronic ground state for fixed  $R$ .
- $F_v^{1\sigma_g}(R)$  is a vibrational level for the nuclear motion. [We show results for  $v = 0$ , but can take arbitrary distributions.]
- The screening parameters  $q_n = 0.03$  and  $q_e = 1.0$  are used to smooth out the Coulomb singularity.
- The term  $z f(t) \mathcal{E}_0 \sin(\omega t)$  represents the effect of the electric field with amplitude  $\mathcal{E}_0$  and angular frequency  $\omega$ .
- Finally,  $f(t)$  is a smooth turn-on/turn-off function for the field.
- We have used intensities of  $1-2 \times 10^{18} \text{ W/m}^2$ , turn-on/turn-off times of **2-4 periods**, and on-times of about **10 periods**.
- We also **follow the system for several periods after the pulse is over**.

## Extracting Physics

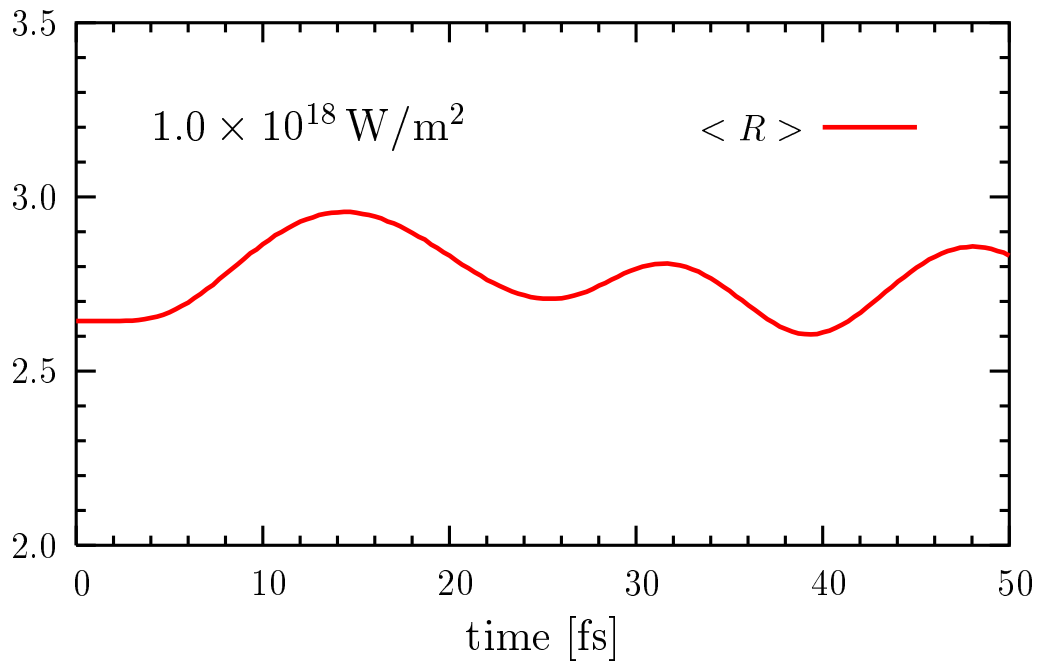
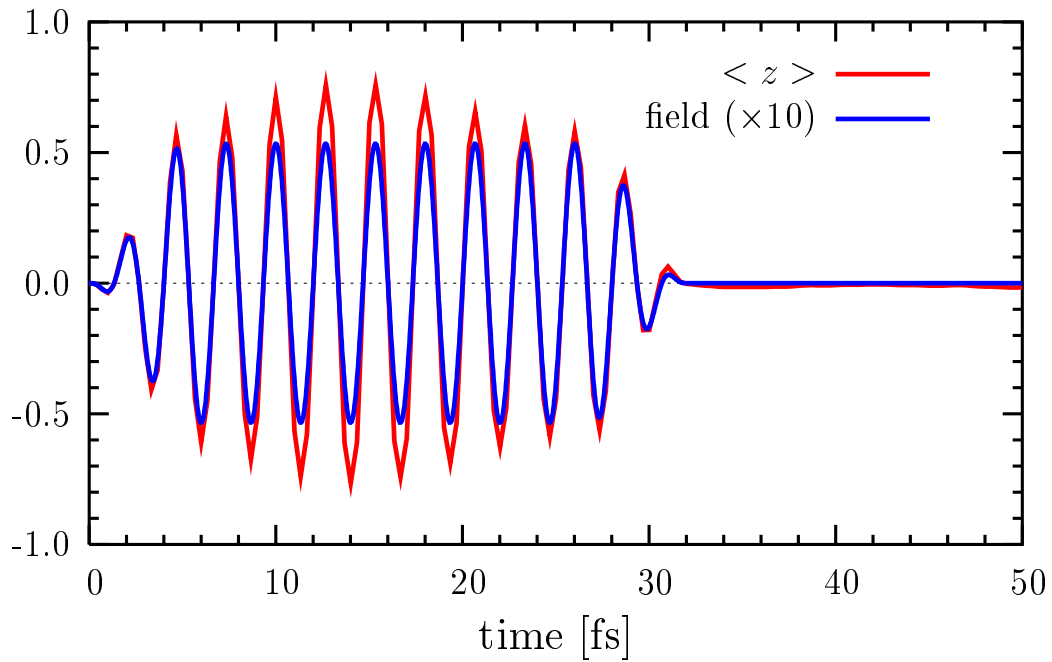
- Standard observables to calculate are  $\langle R(t) \rangle$  and  $\langle z(t) \rangle$ .
- However,  $\langle R(t) \rangle$  has limited meaning since it includes vibration, dissociation, and Coulomb explosion; the same is true for  $\langle z(t) \rangle$ .
- Looking at the complete nuclear and electronic distributions reveals significantly more information.
- In the following movies, watch for:
  - bond-stretching and bond-healing
  - vibrational excitation
  - dissociation (electron stays with one of the protons)
  - ionization (electron leaves the system, so protons undergo “Coulomb explosion”)
- Ultimately, we plan to project to particular final states to find, for example, the probability of exciting a particular vibrational and/or electronic state. In general, such a probability is given by

$$P_v^{n\sigma_{g,u}}(t) \equiv \left| \int_0^\infty dR \int_\infty^\infty dz F_v^{n\sigma_{g,u}}(R) G_R^{n\sigma_{g,u}}(z) \Psi(R, z, t) \right|^2$$

- To get a cross section, we need to take  $\lim_{t \rightarrow \infty} P_v^{n\sigma_{g,u}}(t)$  and introduce a conversion factor.
- Similarly, we plan to determine ionization as well as dissociation cross sections.

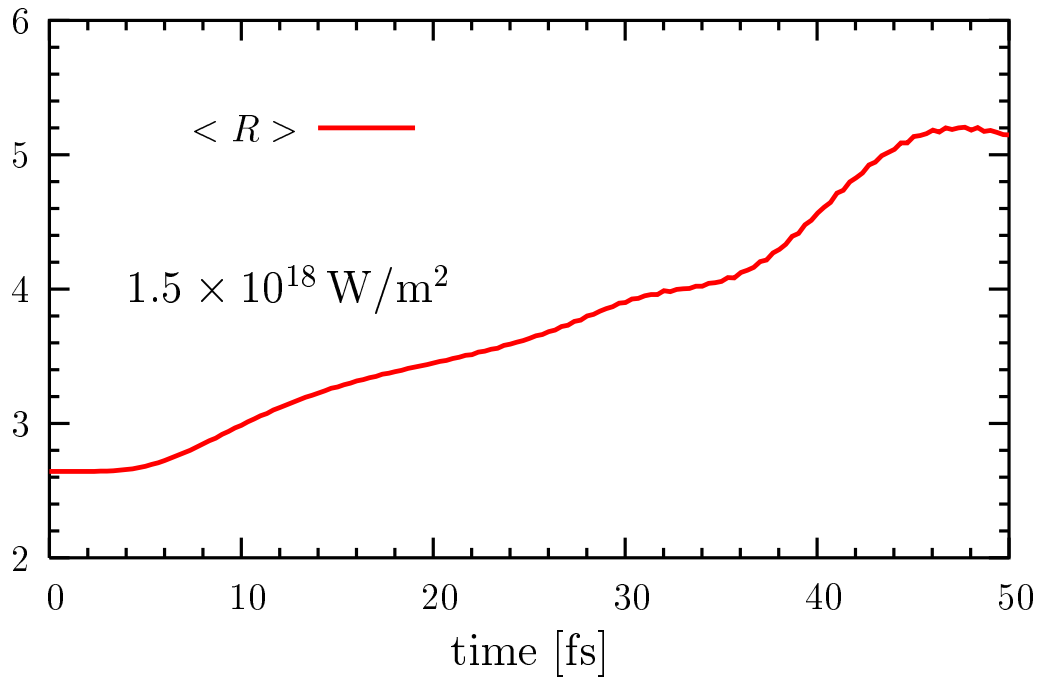
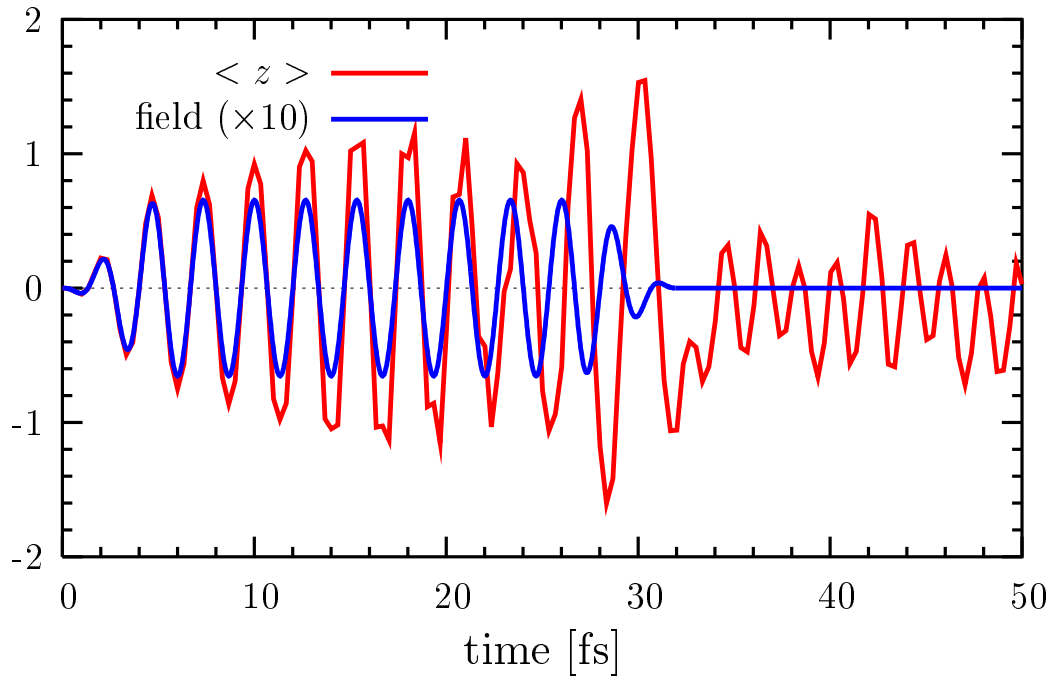
## Expectation Values

turn-on/off: 2 periods; on-time: 8 periods; time after: 8 periods



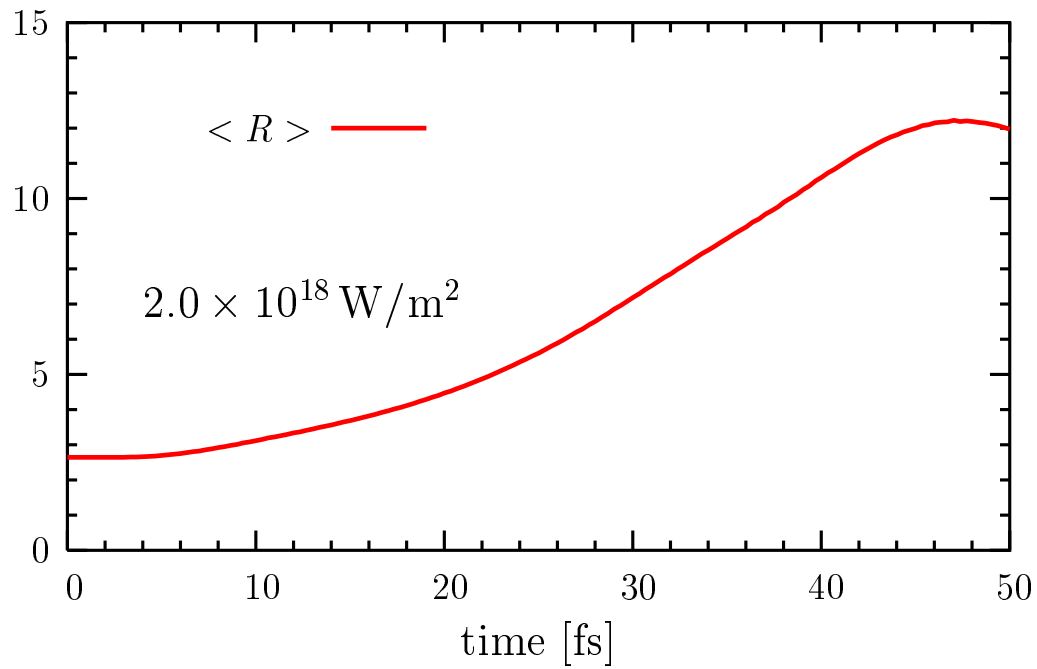
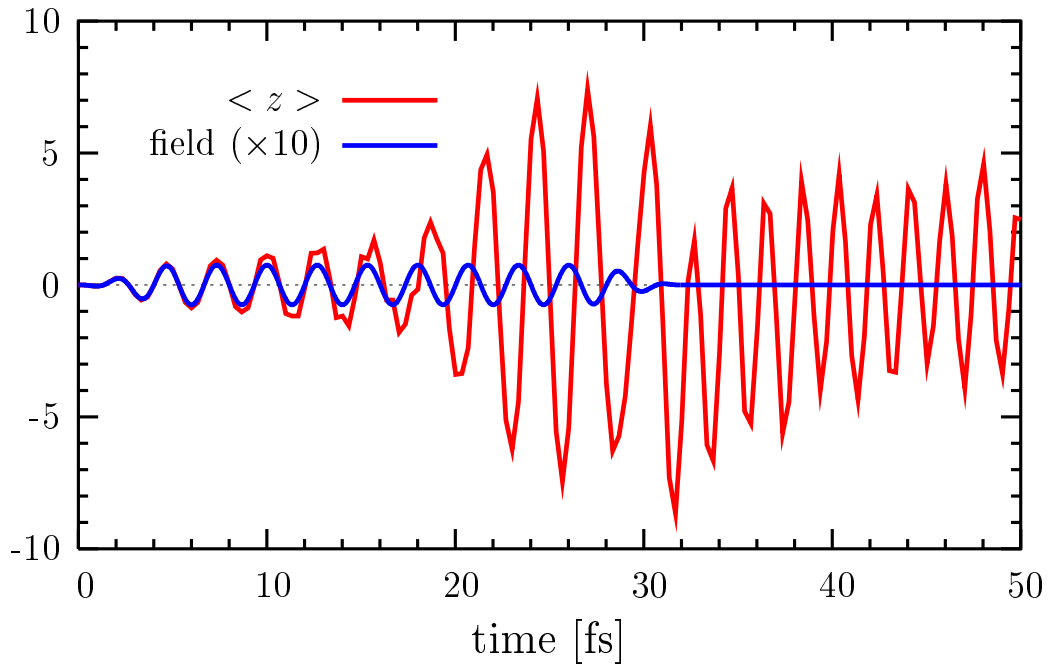
## Expectation Values

turn-on/off: 2 periods; on-time: 8 periods; time after: 8 periods



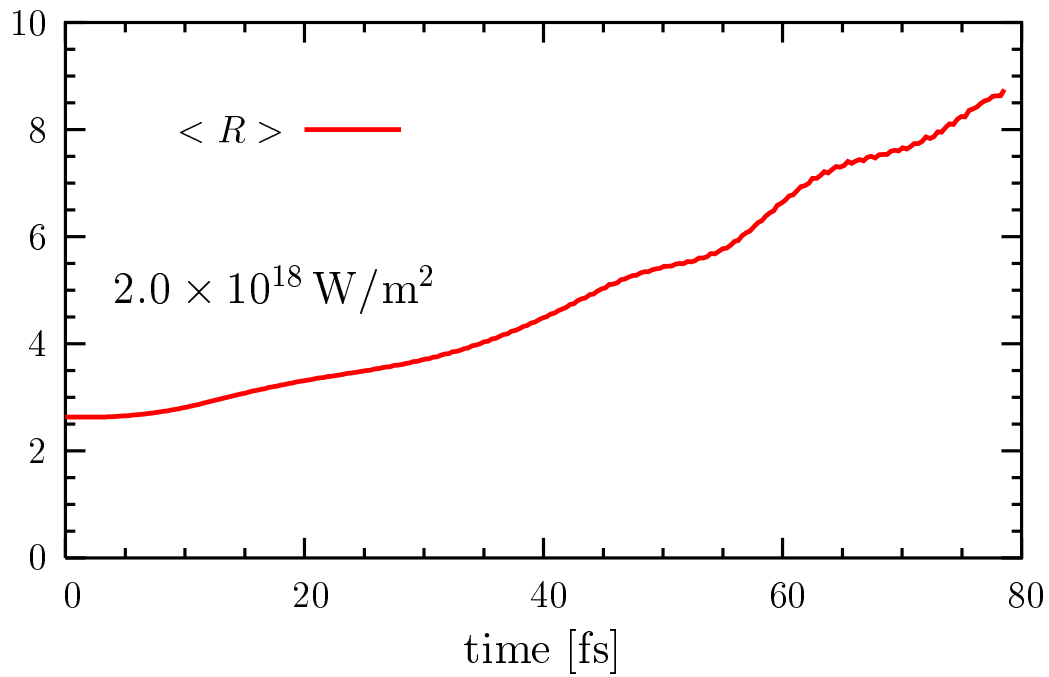
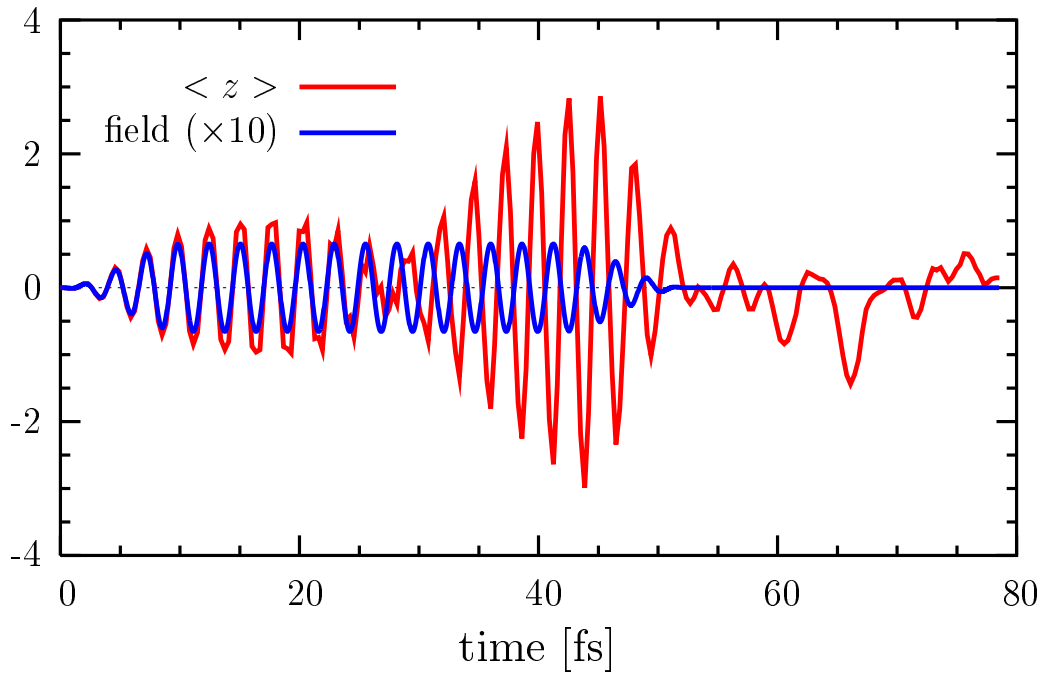
## Expectation Values

turn-on/off: 2 periods; on-time: 8 periods; time after: 8 periods



## Expectation Values

turn-on/off: 4 periods; on-time: 12 periods; time after: 10 periods





## Conclusions and Outlook

- Our method has been checked with good agreement against benchmark results for the Temkin-Poet model of e-H collisions [Phys. Rev. A 65 (2002) 060701(R)] and the  $\text{H}_2^+$  1D-problem.
- By avoiding boundary masking, we can extract useful physics about the final state of the system using simple projection methods.
- By projecting onto different states, we can separately find probabilities for excitation, ionization, and dissociation processes.
- We note the importance of watching the evolution of the system after the laser pulse is over ( $\rightarrow$  “bond healing”).
- The Temkin-Poet code has been parallelized. We are in the process of doing the same for the  $\text{H}_2^+$  problem, which will allow us to use even bigger grids and longer propagation times.
- Possibilities to extend the analysis beyond one dimension with two degrees of freedom will be explored.
- Movies have been created to enhance the visualization.